

## 6.7 Non-homogeneous equations; introduction to Green's functions

### 6.7.1 Non homogeneous (regular) S.L. problems (ODEs)

Given the ODE  $\frac{1}{r(x)} [(p(x)u)'] + q(x)u = F(x)$

$$\text{with bcs } \begin{cases} \alpha u(a) + \beta u'(a) = 0 \\ \gamma u(b) + \delta u'(b) = 0 \end{cases}$$

① Seek solutions of the homogeneous eigenvalue eq.

$$\frac{1}{r(x)} [(p(x)u)'] + q(x)u = -\lambda u$$

→ this yields the eigenfunctions  $\{v_n\}$  and eigenvalues  $\{\lambda_n\}$

② Write  $F(x) = \sum_n b_n v_n(x)$

(with  $b_n = \int_a^b r(x) F(x') v_n(x') dx'$ , if the  $v_n$ s are properly normalized)

Then since we know that the solution can also be written as

$$u(x) = \sum_n a_n v_n(x)$$

we can write

$$\frac{1}{r(x)} [(p(x)u)'] + q(x)u = \sum_n -\lambda_n a_n v_n(x) = \sum_n b_n v_n(x)$$

and by identification,  $a_n = -\frac{b_n}{\lambda_n}$

$$\begin{aligned} \Rightarrow u(x) &= \sum_n -\frac{b_n}{\lambda_n} v_n(x) = -\sum_n \int_a^b \frac{r(x') F(x') v_n(x') v_n(x) dx'}{\lambda_n} \\ &= \int_a^b G(x, x') F(x') dx' \end{aligned}$$

where  $G(x; x') = -\sum_n \frac{1}{\lambda_n} v_n(x') v_n(x) r(x')$

- $G(x; x')$  is called the Green's function of the S. L. problem
- It only depends on the characteristics of the homogeneous problem ( $\{v_n\}, \{\lambda_n\}$ ) but, when integrated through with the forcing term  $F(x)$ , yields the solution of the forced problem
- Note that if the  $\{v_n\}$  are not normalized then

$$G(x; x') = -\sum_n \frac{1}{\|v_n\|^2} \frac{r(x')}{\lambda_n} v_n(x') v_n(x)$$

$$\text{where } \|v_n\|^2 = \int_a^b r(x) v_n(x)^2 dx$$

Example Consider

$$y'' + y = 3 \sin(2\pi x) \quad \begin{array}{l} y(0) = 0 \\ y(1) = 0 \end{array}$$

We seek the eigenfunctions of  $y'' + y = -\lambda y$

$$\rightarrow y'' + (1+\lambda)y = 0$$

$$\text{so } y = \alpha \cos(\sqrt{1+\lambda}x) + \beta \sin(\sqrt{1+\lambda}x)$$

$$\text{with } \begin{cases} \alpha = 0 \\ \sqrt{1+\lambda_n} = n\pi \end{cases}$$

$$\Rightarrow \begin{cases} \lambda_n = n^2\pi^2 - 1 \\ v_n(x) = \sin(n\pi x) \end{cases}$$

$$\begin{aligned} \Rightarrow \text{The Green's function } G(x, x') &= \sum_n \frac{\sin(n\pi x) \sin(n\pi x')}{\lambda_n \|\sin(n\pi x)\|^2} \\ &= \sum_n \frac{2}{n^2\pi^2 - 1} \sin(n\pi x) \sin(n\pi x') \end{aligned}$$

so that

$$\sum_n \dot{a}_n(t) v_n(x) + \sum_n \lambda_n a_n(t) v_n(x) = \sum_n b_n(t) v_n(x)$$

and (by orthogonality):

$$\dot{a}_n + \lambda_n a_n = b_n(t)$$

⇒ integrating factor method:

$$\frac{d}{dt} (a_n e^{\lambda_n t}) = b_n(t) e^{\lambda_n t}$$

$$\text{so } a_n(t) e^{\lambda_n t} - a_n(0) = \int_0^t b_n(t') e^{\lambda_n t'} dt'$$

$$\Rightarrow a_n(t) = a_n(0) e^{-\lambda_n t} + e^{-\lambda_n t} \int_0^t b_n(t') e^{\lambda_n t'} dt'$$

Putting it all together we find that

$$\begin{aligned} u(x,t) &= \sum_n a_n(0) e^{-\lambda_n t} v_n(x) + \int_0^t \sum_{n'} e^{-\lambda_{n'}(t-t')} v_{n'}(x) b_{n'}(t') dt' \\ &= \sum_n a_n(0) e^{-\lambda_n t} v_n(x) + \int_0^t \int_a^b \sum_n e^{-\lambda_n(t-t')} v_n(x) v_n(x') r(x') F(x',t') dx' dt' \end{aligned}$$

So we can write

$$u(x,t) = \sum_n a_n(0) e^{-\lambda_n t} v_n(x) + \int_0^t \int_a^b G(x,t; x',t') F(x',t') dx' dt'$$

$$\text{with } G(x,t; x',t') = \sum_n e^{-\lambda_n(t-t')} v_n(x) v_n(x') r(x')$$

Here  $G$  is called the Heat Equation kernel.  
↳ another example of a Green's function.

↳  $u$  is the sum of

- the solution to the problem with no forcing
- the weighted integral of  $F(x,t)$  with the Green's function.

so the solution to the problem is  $y(x) = \int_0^1 G(x, x') F(x') dx'$

$$y(x) = \int_0^1 \sum_n \frac{3 \sin(2\pi x')}{1 - n^2 \pi^2} 2 \sin(n\pi x) \sin(n\pi x') dx'$$
$$= \frac{3}{1 - 4\pi^2} \sin(2\pi x)$$

### 6.7.2 Application to parabolic/hyperbolic PDEs

Now consider either  $u_t - \frac{1}{r(x)} [(p(x)u')' + q(x)u] = F(x,t)$

or  $u_{tt} - \frac{1}{r(x)} [(p(x)u')' + q(x)u] = F(x,t).$

Idea: Solve the associated Sturm-Liouville problem

$$\frac{1}{r(x)} [(p(x)u')' + q(x)u] + \lambda u = 0$$

to find the eigenvalues and eigenfunctions  
 $\{v_n\}, \{\lambda_n\}$

then expand

$$F(x,t) = \sum_n b_n(t) v_n(x)$$

(in this case,  $b_n(t) = \int_a^b F(x,t) r(x) v_n(x) dx$ )

Assume a solution of the form

$$u(x,t) = \sum_n a_n(t) v_n(x)$$

and try the ansatz into the equation:

Parabolic case  $\therefore \sum_n \dot{a}_n(t) v_n(x) - \frac{1}{r(x)} \left[ \left( p(x) \sum_n a_n(t) v_n'(x) \right)' + q(x) \sum_n a_n(t) v_n(x) \right]$

$$= \sum_n b_n(t) v_n(x)$$

## Example of the drunks exiting the pub.

Recall:

$$\frac{\partial p}{\partial t} = k \frac{\partial^2 p}{\partial x^2} + S(x, t)$$

$$\begin{cases} p(x, 0) = 0 \\ \frac{\partial p}{\partial x} = 0 \text{ at } x = 0, L \\ S(x, t) = S_0 e^{-\frac{t}{\tau}} \delta(x - \frac{L}{2}) \text{ for } t > 0 \end{cases}$$

(take  $\tau = 0$ ).

Homogeneous problem; separation of variables to get spatial eigenmodes  $\Rightarrow$

$$\begin{cases} v_n(x) = \cos\left(\frac{n\pi x}{L}\right) \\ \lambda_n = \frac{n^2 \pi^2}{L^2} \end{cases}$$

So, by the previous calculation, we have

$$u(x, t) = \sum_{n=0}^{\infty} a_n(t) e^{-\lambda_n t} v_n(x) + \int_0^t \int_0^L S(x', t') G(x, x'; t, t') dx' dt'$$

where  $a_n(t)$  are obtained by fitting  $u$  to initial conditions

$$u(x, 0) = \sum_{n=0}^{\infty} a_n(0) v_n(x) = 0 \Rightarrow a_n(0) = 0$$

and where  $G(x, x'; t, t') = \sum_{n=0}^{\infty} e^{-\lambda_n(t-t')} \frac{v_n(x') v_n(x) r(x)}{\|v_n\|^2}$

$$\begin{aligned} \Rightarrow u(x, t) &= \int_0^t \int_0^L \sum_{n=0}^{\infty} S_0 e^{-\frac{t'}{\tau}} \delta(x' - \frac{L}{2}) e^{-\frac{n^2 \pi^2}{L^2}(t-t')} \frac{1}{\|v_n\|^2} \cos\left(\frac{n\pi x'}{L}\right) \cos\left(\frac{n\pi x}{L}\right) dx' dt' \\ &= \int_0^t \sum_{n=0}^{\infty} S_0 e^{-\frac{t'}{\tau}} e^{-\frac{n^2 \pi^2}{L^2}(t-t')} \cos\left(\frac{n\pi}{2}\right) \cos\left(\frac{n\pi x}{L}\right) dt' \cdot \frac{1}{\|v_n\|^2} \\ &= \sum_{n=0}^{\infty} S_0 \cos\left(\frac{n\pi}{2}\right) \cos\left(\frac{n\pi x}{L}\right) \frac{1}{\frac{n^2 \pi^2}{L^2} - \frac{1}{\tau}} \left[ e^{-\frac{t}{\tau}} - e^{-\frac{n^2 \pi^2}{L^2} t} \right] \frac{1}{\|v_n\|^2} \end{aligned}$$